

Theoretical determination of the fractal dimension of fluid parcel trajectories in large and meso-scale flows

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Recent work suggests that Lagrangian parcel trajectories in large and meso-scale flows can have a fractal dimension (OSBORNE, et al, 1986, 1989; PROVENZALE, et al, 1989). Drifter trajectories may be viewed as fractal curves in the plane of the ocean's surface and have typical fractal dimensions of about 1.30 ± 0.06 in a range of scales normally attributed to geostrophic turbulence. There are several fundamental questions which we address with regard to these results (OSBORNE and PROVENZALE, 1989; OSBORNE and CAPONIO, 1990a, b):

- (1) One normally thinks of solutions to partial differential equations as being reasonably smooth, well-behaved differentiable functions. How can it theoretically be possible for the *dynamical* motions of particle trajectories to *physically* be fractal curves?
- (2) What are the *physical implications* of fractal trajectories on fluid flows in general?
- (3) What new, *unique physical information* can the fractal dimension give us with regard to large scale flows?
- (4) Can the fractal dimension be of use in solving the *inverse problem*, i.e. determination of the general circulation from drifter trajectories?
- (5) What relation does the fractal dimension have to fluid properties such as *anisotropic and anomalous diffusion* in fluids?
- (6) What are the implications on *fractal front propagation* and on *fractal frontogenesis*?

To address these questions we discuss a simple, nonlinear Hamiltonian model for describing particle motions in large and meso-scale oceanic and atmospheric flows. The model predicts nonlinear fluid parcel trajectories in 2-D where the stream function power spectrum has the form $k^{-\gamma}$ (γ const), while the velocity spectrum is given by $k^{-\delta}$, $\delta = \gamma - 3$. The equations of motion are

$$\dot{\mathbf{x}} = u_{rms} \hat{\mathbf{k}} \times \nabla \psi(\mathbf{x}, t)$$

The over-dot denotes time derivative, $\mathbf{x}(t) = [x(t), y(t)]$ is the parcel position in the xy-plane, $\hat{\mathbf{k}}$ is the unit vector perpendicular to this plane, $\nabla = (\partial_x, \partial_y, \partial_z)$ and u_{rms} is the assumed rms velocity of the flow. $\psi(\mathbf{x}, t)$ is the stream function or Hamiltonian, which is here given a stochastic representation in two dimensions. The dispersion relation is $\omega = uk$, where $k = |k|$, and u is constant. Using this model as a basis we apply both analytical and numerical methods to show that fluid particle trajectories: (1) while differentiable and chaotic at small scale, may be *fractal space curves* at larger scale and (2) undergo *anomalous transport*.

We identify a "stochasticity parameter" $\mu = u/u_{rms}$ which characterizes the flow: (a) When $\mu = 0$ the Hamiltonian is time independent and exactly integrable, but the system is nevertheless fully nonlinear. For specific initial conditions $\mathbf{x}(0) = \mathbf{x}_0$, $\mathbf{x}(t)$ lies on the contour given by $\psi(\mathbf{x}_0) = \text{const}$. For $0 \leq \delta \leq 1$, $\mathbf{x}(t)$ is a fractal curve in the xy-plane with dimension $D = 2/(\delta + 1)$, $1 \leq D \leq 2$, and has *anomalous absolute diffusion* $\langle |\mathbf{x}(t) - \mathbf{x}(0)|^2 \rangle = D t^{2/D}$, for D the diffusion coefficient. (b) For $\mu \ll 1$ and finite, the flow is chaotic and of *KAM* (Kolmogorov-Arnold-Moser) type, i.e. the Hamiltonian consists of $\psi(\mathbf{x})$ plus a perturbation. The *KAM "surfaces"* are very nearly fractal contours of $\psi(\mathbf{x})$. $\psi(\mathbf{x}, t)$ varies slowly in time and stochastic layers are formed. (c) For $\mu < 1$, the flow is chaotic and lies in a *trapping regime* characterized by *vortices* and *vortex hopping*. (d) When $\mu \approx 1$, clustering occurs, and isolated clumps of activity appear, evidently the remnants of the vortices. (e) For $\mu > 1$, the flow is *fully stochastic*, equivalent to a *nongaussian random walk* in the plane. Example trajectories are given in Figure 1.

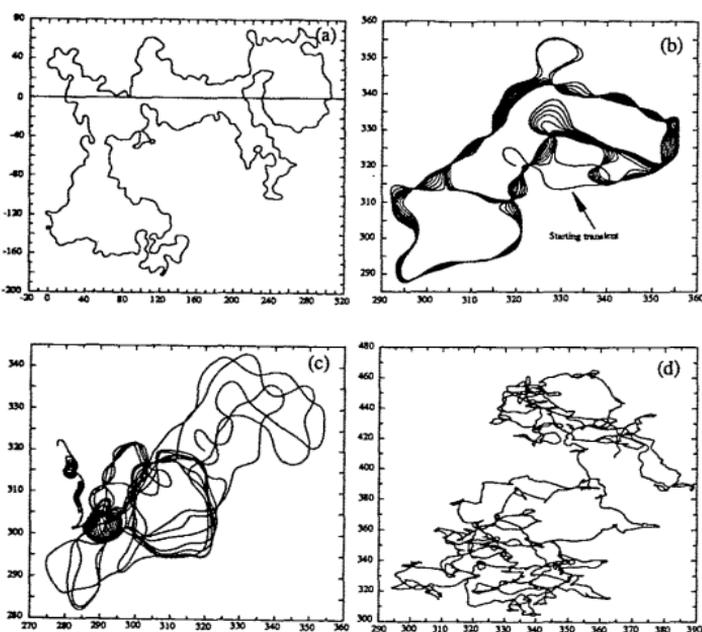


Figure 1. Particle trajectories for a velocity spectrum $k^{-0.5}$; the units are kilometers; $u_{rms} = 7$ km/day. (a) $\mu = 0$, the motion is periodic, and the trajectory is a fractal curve with dimension $D = 1.33$. (b) $\mu = 0.0014$, the motion is of KAM type and stochastic layers form. (c) $\mu = 0.014$, the motion lies in the vortex/vortex hopping regime. (d) $\mu = 1.4$, the motion is fully stochastic, a nongaussian random walk in the plane.

REFERENCES

- OSBORNE, A. R., KIRWAN, JR., A. D., PROVENZALE, A. and BERGAMASCO, L., 1986 A search for chaotic behavior in large and mesoscale motions in the Pacific Ocean. *Physica D* 23: 75-83.
- OSBORNE, A. R. and PROVENZALE, A., 1989 Finite correlation dimension for stochastic systems with power-law spectra, *Physica D* 35: 357-381.
- OSBORNE, A. R., KIRWAN, JR., A. D., PROVENZALE, A. and BERGAMASCO, L., 1989 Fractal drifter trajectories in the Kuroshio extension, *Tellus* 41A: 416-435.
- PROVENZALE, A., OSBORNE, A. R., KIRWAN, JR., A. D. and BERGAMASCO, L., 1990 The study of fluid parcel trajectories in large scale ocean flows. In *Nonlinear Topics in Ocean Physics*, A. R. Osborne ed., Elsevier, Amsterdam.
- OSBORNE, A. R., and CAPONIO, R., 1990 Fractal Trajectories and Anomalous Diffusion for Chaotic Particle Motions in 2-D Turbulence, submitted for publication.
- OSBORNE, A. R., and CAPONIO, R. (1990) The Transition From Chaos to Stochasticity in 2-D Turbulence, In *Nonlinear and Turbulent Processes in Physics*, A. G. Sitenko, V. E. Zakharov and V. M. Chernoussenko eds.